## E-field Systematic in a Distribution

Because of the finite curvature of the quadrupoles, the E-field correction to  $\omega_a$  can depend on the betatron amplitude as well as the momentum. The muon that circulates at the magic radius with zero betatron oscillation amplitude will see zero electric field and no correction is required. But the muon with that same momentum and finite betatron amplitude will oscillate about the magic radius and its precession will indeed be effected by the E-field. If the quad field were strictly antisymmetric about the magic radius and the betatron trajectory strictly symmetric, the net contribution would be zero. But neither condition is true. The contribution from the betatron amplitude was estimated analytically in doc-db 12575-v2. Here we explore amplitude dependence in simulation.

## **Integrated Electric Field**

The dependence of  $\omega_a$  on electric field is given by

$$\vec{\omega}_a = -\frac{q}{m} \left[ a_\mu \vec{B} - \left( a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right] \tag{1}$$

A electric field along the trajectory of a muon will vary as it executes betatron oscillations and the contribution of the field to  $\vec{\omega}_a$  can be characterized in terms of an average.

$$\vec{\omega}_a = -\frac{q}{m} \left[ a_\mu \vec{B} - \left( a_\mu - \frac{1}{\gamma^2 - 1} \right) \int \frac{\vec{\beta} \times \vec{E}}{c} \frac{dt}{T} \right]$$
 (2)

The fractional correction to  $\vec{\omega}_a$  is

$$\vec{C}_e \sim 2 \frac{\Delta p}{p} \int \frac{\vec{\beta} \times \vec{E}}{Bc} \frac{dt}{T}$$

Of course the component of  $\vec{C}_e$  in the vertical is the largest.

In simulation we integrated  $\vec{\beta} \times \vec{E}$  along the trajectory of each muon to estimate  $C_e$ . Results for individual trajectories are shown in Figure 1 where the average integrated  $\vec{\beta} \times \vec{E}$  is plotted as a function of time for a muon with  $\Delta p/p = 0.2\%$  and betatron amplitude  $a_x = 0.74, 39.7, 64.3$ , and 97.7 mm-mrad respectively. (The amplitude  $a_x = (\gamma x^2 + 2\alpha x x' + \beta (x')^2)^{1/2}$  The effect of the electric field is mitigated (a bit) with increasing betatron amplitude. The oscillation of the correction at early times reflects the alternating sign of the contribution with radial displacement with a maximum discrepancy at the level of  $\sim 5\%$ .

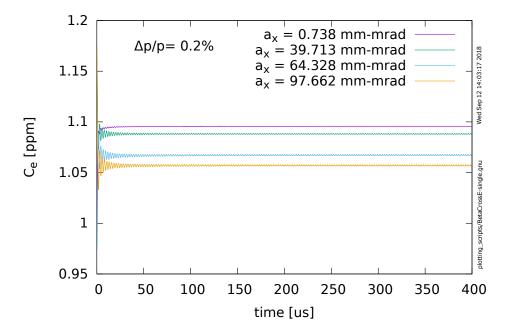


Figure 1: Efield correction to  $a_{\mu}$  at  $\Delta p/p = 0.2\%$  and 4 different betatron amplitudes

## Distribution

In the same way we evaluate the electric field correction for a distribution of particles. The initial distribution is generated in the start to ring simulation[1]. The temporal distribution is as measured[2]. The particles are propagated through the injection line, inflector and into the ring. The kicker pulse shape is assumed as measured but with design field of 292 G. The captured muons (about 2% of the initial distribution) are propagated around the ring through 4000 turns. There is no decay. The radial phase space distribution at turn 4000 is shown in Figure 2 and the momentum distribution in Figure 3.

The efield correction is shown in Figure 4 as a function of momentum offset. The quadratic dependence is evident. The correction is plotted as a function of betatron amplitude in Figure 5. There is no obvious correlation..

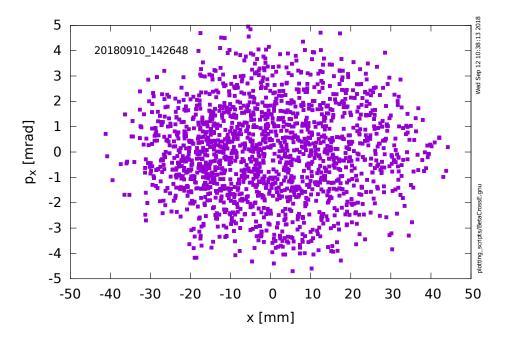


Figure 2: Radial phase space at turn 4000

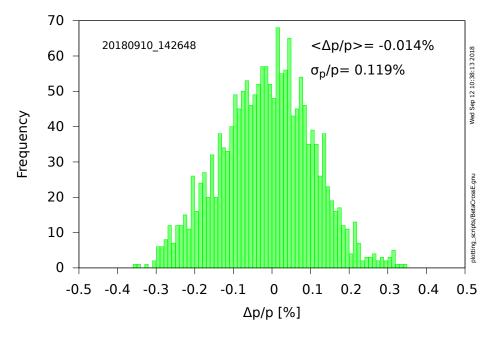


Figure 3: Momentum distribution

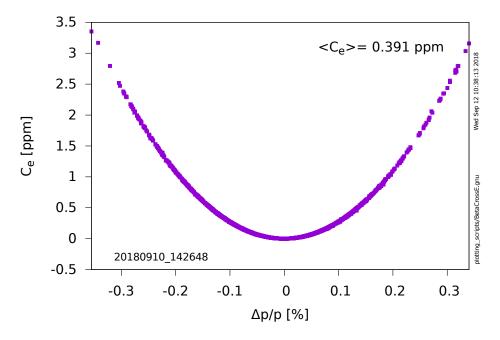


Figure 4: Efield correction as a function of momentum offset

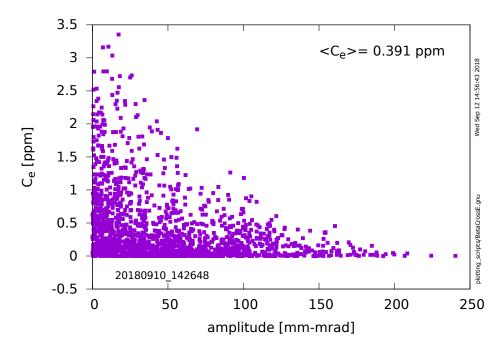


Figure 5: Efield correction vs betatron amplitude

## Bibliography

- $[1]\,$  D. Stratakis, "  $400{,}000$  particles at the end of M5", DocDb 4461-v1
- [2] C. Stoughton, "Simulated pulses for Kicker upgradesâ<br/>Ăİ C. Stoughton in docdb 11721