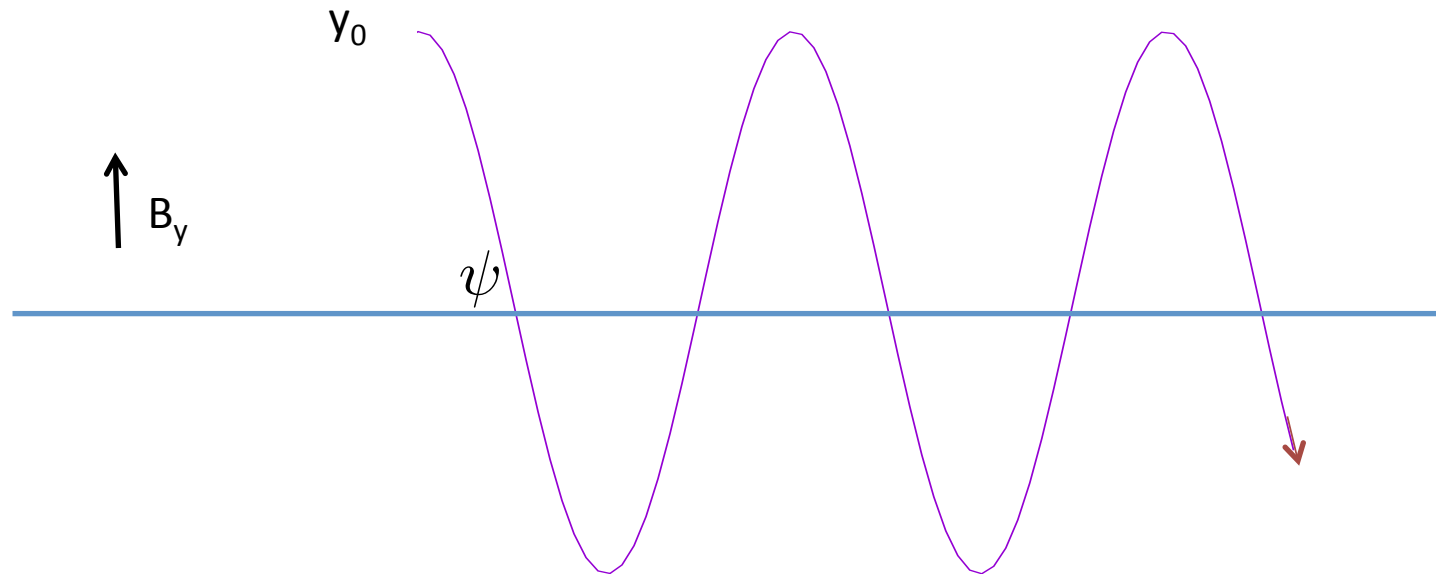


Spin tracking vs Integration and effect of quad nonlinearity

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February 7, 2019

Pitch systematic



$$\frac{\Delta\omega}{\omega} = \langle (1 - \hat{\beta} \times \hat{\mathbf{B}}) \rangle \approx \frac{1}{2} \langle \psi^2 \rangle$$

Pitch correction

3 ways to compute pitch correction in simulation

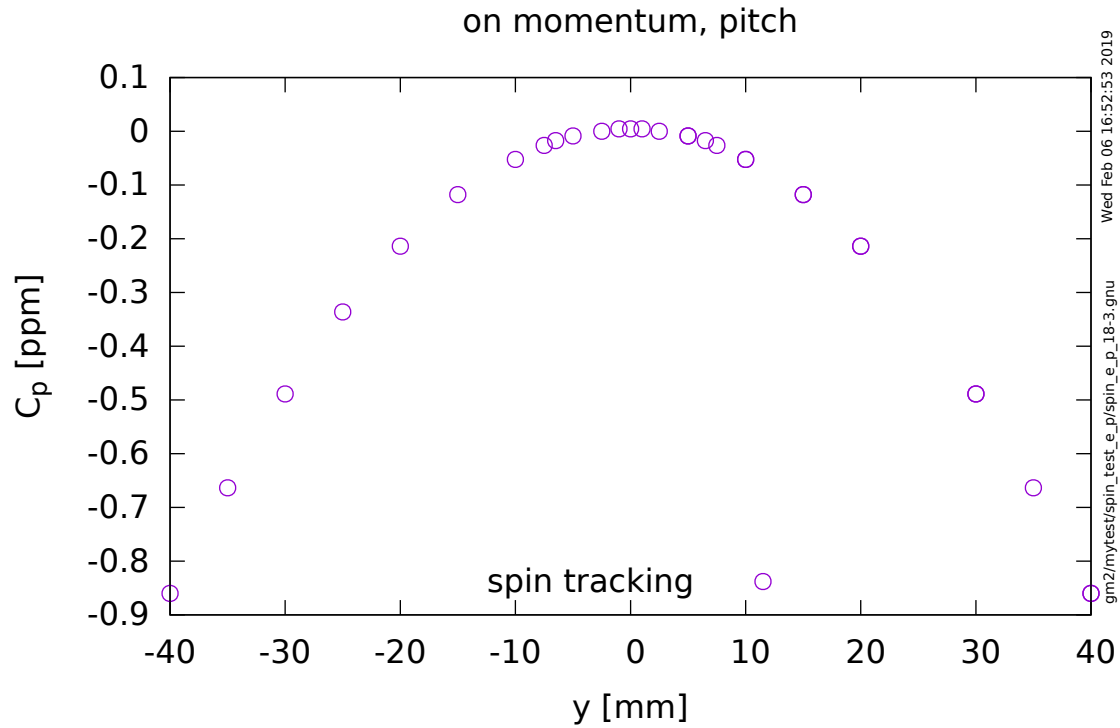
- Spin tracking – *Includes everything, but ppb precision requires many turns*
- Integration along trajectory – *very good approximation far from resonances*

$$C_p = \frac{1}{T} \int_0^T (1 - \hat{\beta} \times \hat{\mathbf{B}}) dt$$

- Measurement of vertical amplitude – *assumes quad linearity*

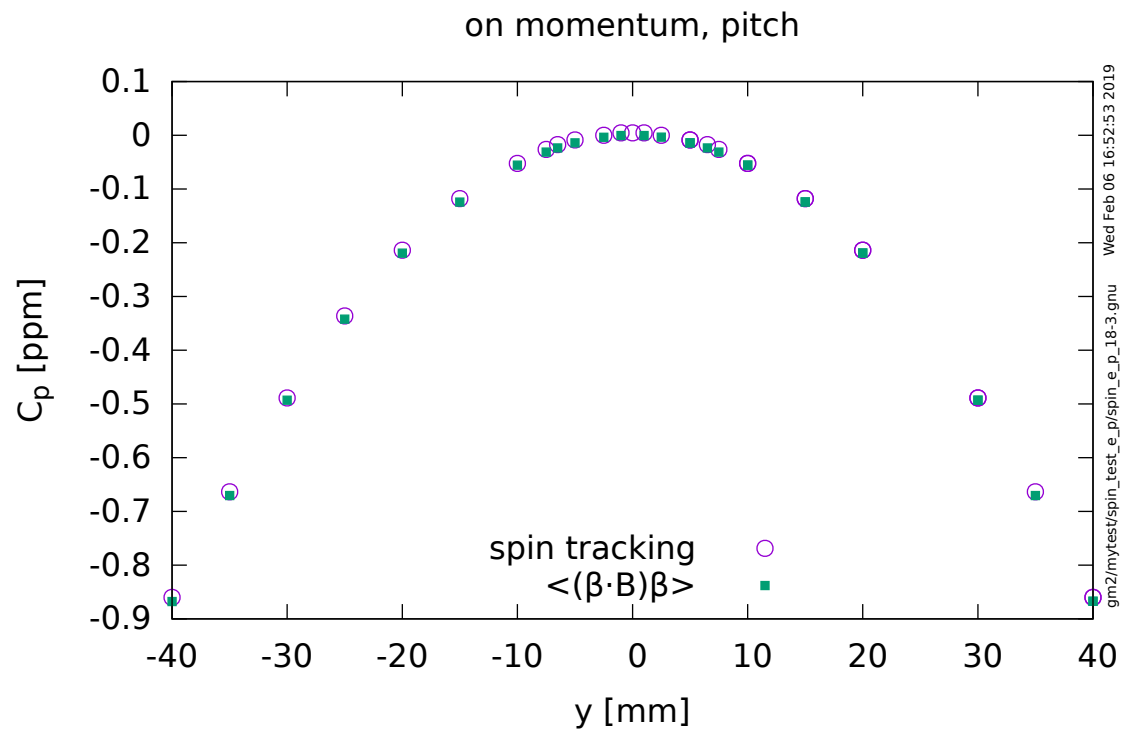
$$C_p = -\frac{n\langle y^2 \rangle}{2R_0^2} = -\frac{\langle y^2 \rangle}{2\beta_y^2} = -\frac{\langle \psi^2 \rangle}{2}$$

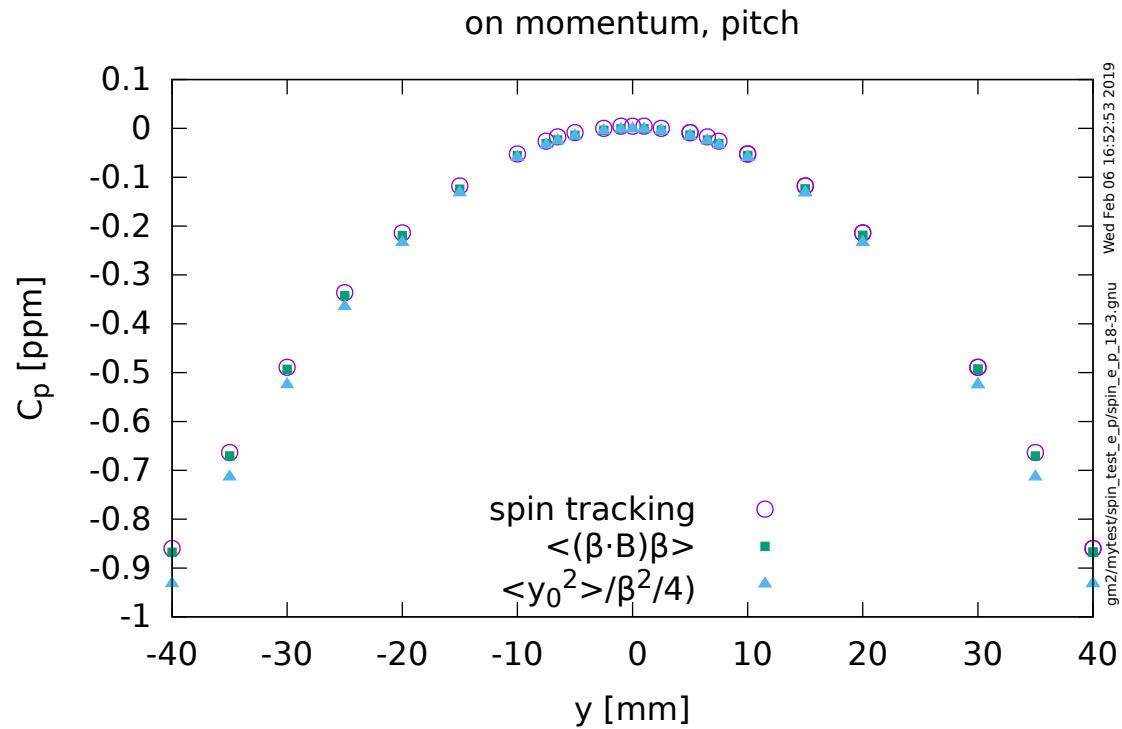
Spin tracking using BMT



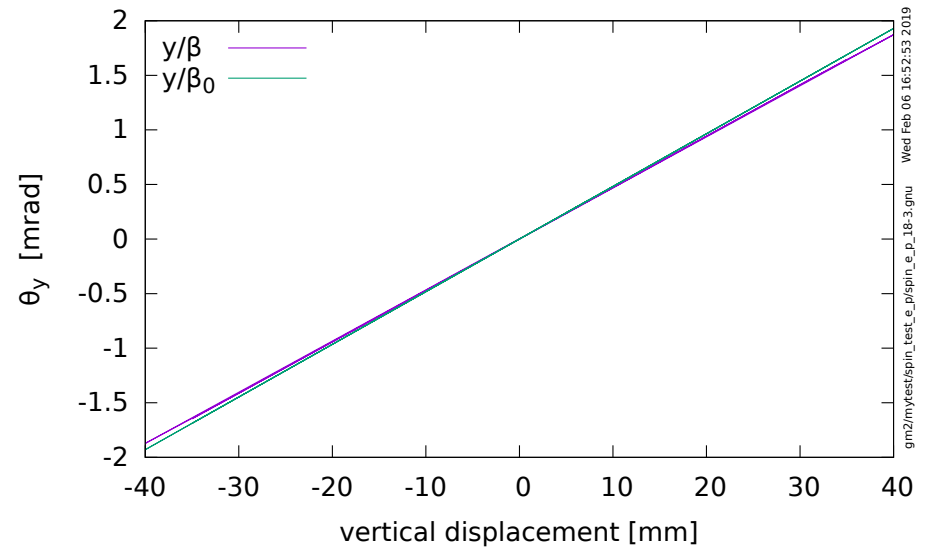
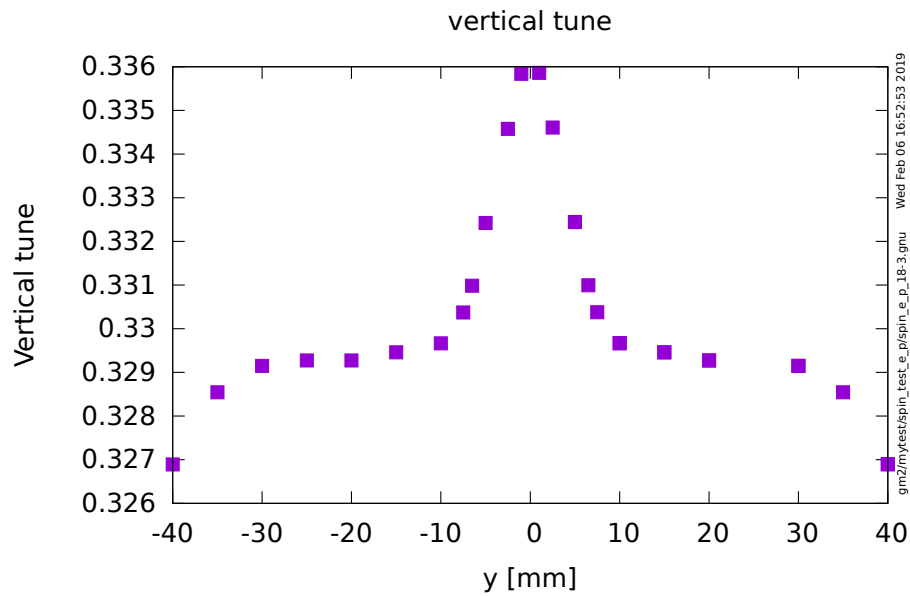
Pitch correction vs vertical amplitude

Spin tracking and 'integration' are in good agreement



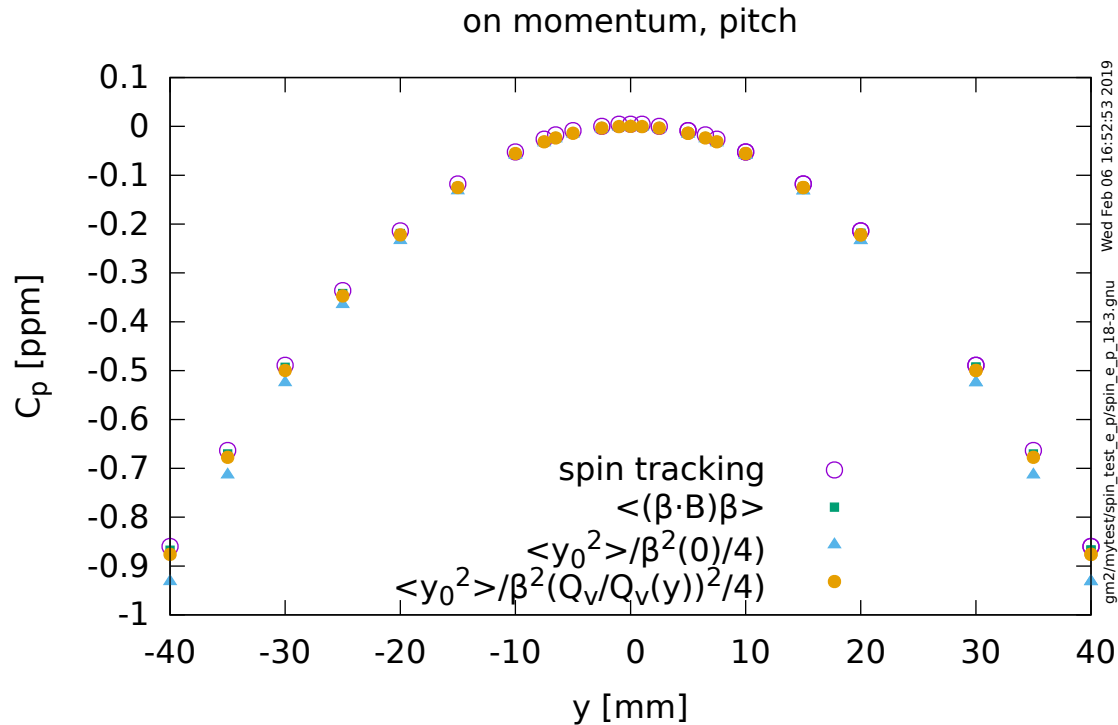


Vertical amplitude $(y_0)/\beta_v$ is not a good measure of angle ψ at large amplitude

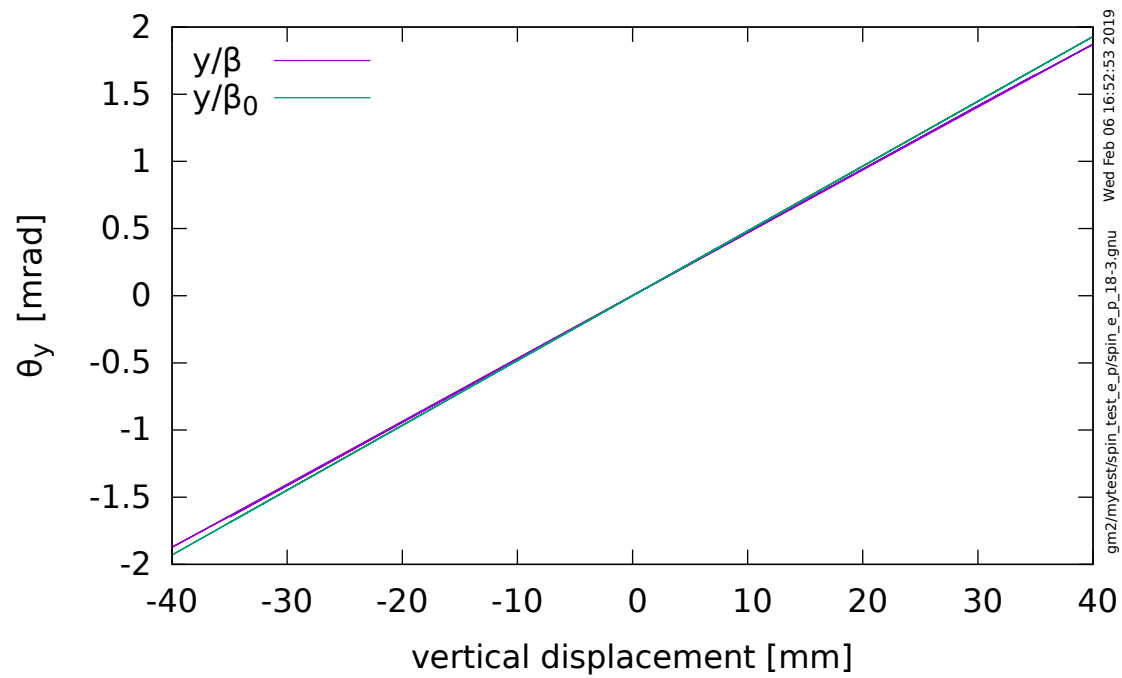


Quad nonlinearity => amplitude dependence of tune and β
 And nonlinear dependence of ψ on y_0

$$\beta(y) = \left(\frac{Q_v(0)}{Q_v(y)} \right)^2 \beta_0$$



- We can correct for the amplitude dependence by measuring the vertical tune
- Alternatively, measure the angular $\langle \psi^2 \rangle$ distribution directly



E field correction

3 ways to compute E-field contribution to ω_a

1. Spin tracking (BMT equation)

2. Integration
$$\vec{C}_e(T) \sim 2 \frac{\Delta p}{p} \frac{1}{T} \int_0^T \frac{\vec{\beta} \times \vec{E}}{Bc} dt$$

a) Integration along trajectory (includes betatron oscillations)

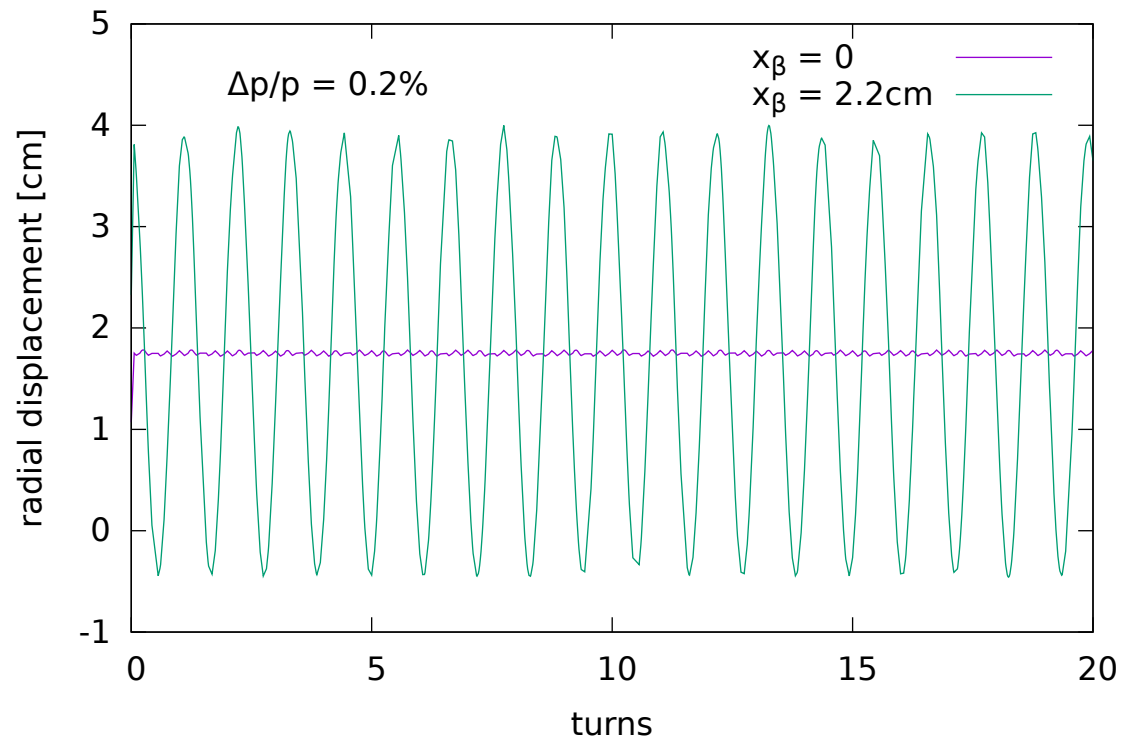
b) Integration along closed orbit ($x = \eta\delta$)

Note that method 2b) is most nearly equivalent to the ‘classic’ method, namely

$$C_E = -2\beta^2 n(1-n) \frac{\langle x_e^2 \rangle}{r_0^2}$$

Compare the 3 methods in simulation to determine

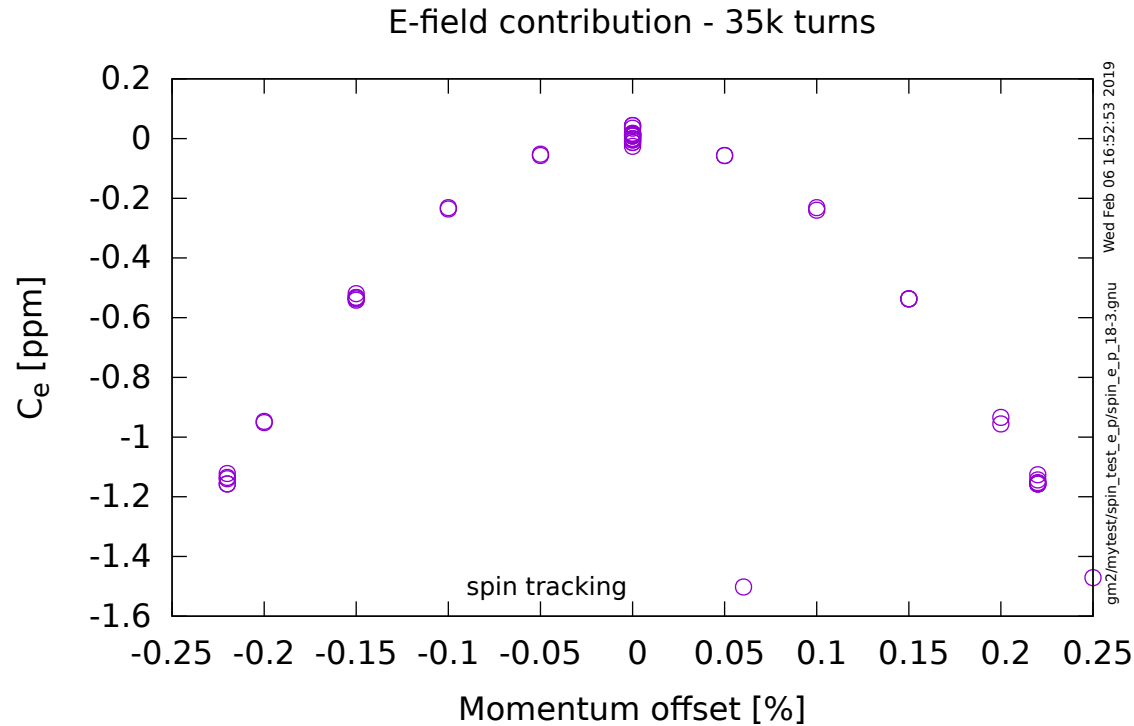
1. If integration is a reliable proxy for spin tracking
2. The size of the contribution from finite betatron oscillation amplitude
3. Effect of quad nonlinearity



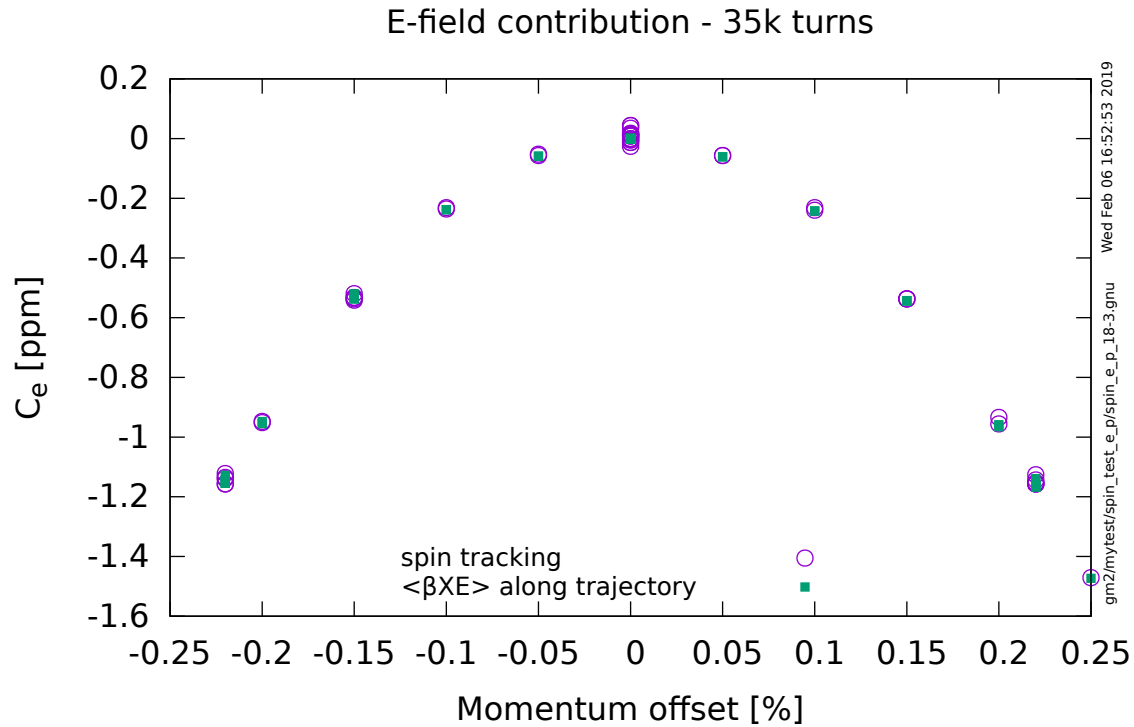
Distinct trajectories with common momentum offset

- For trajectory compute ω_a by spin tracking and by integration

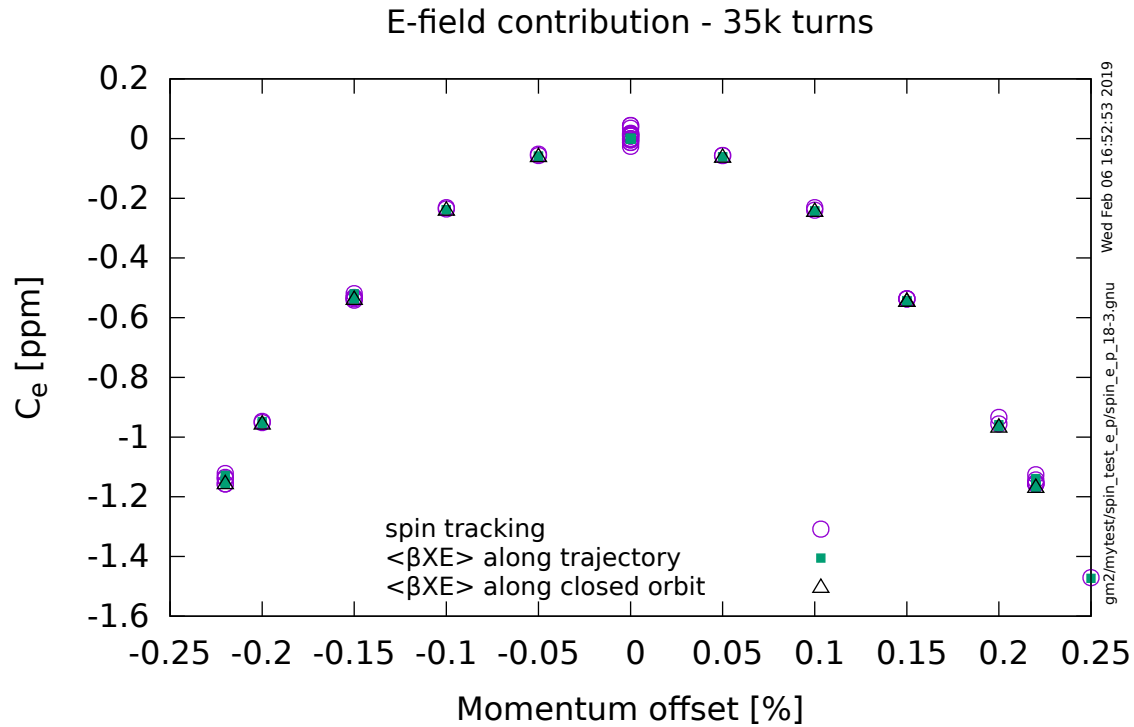
Is the Efield correction independent of the betatron amplitude?



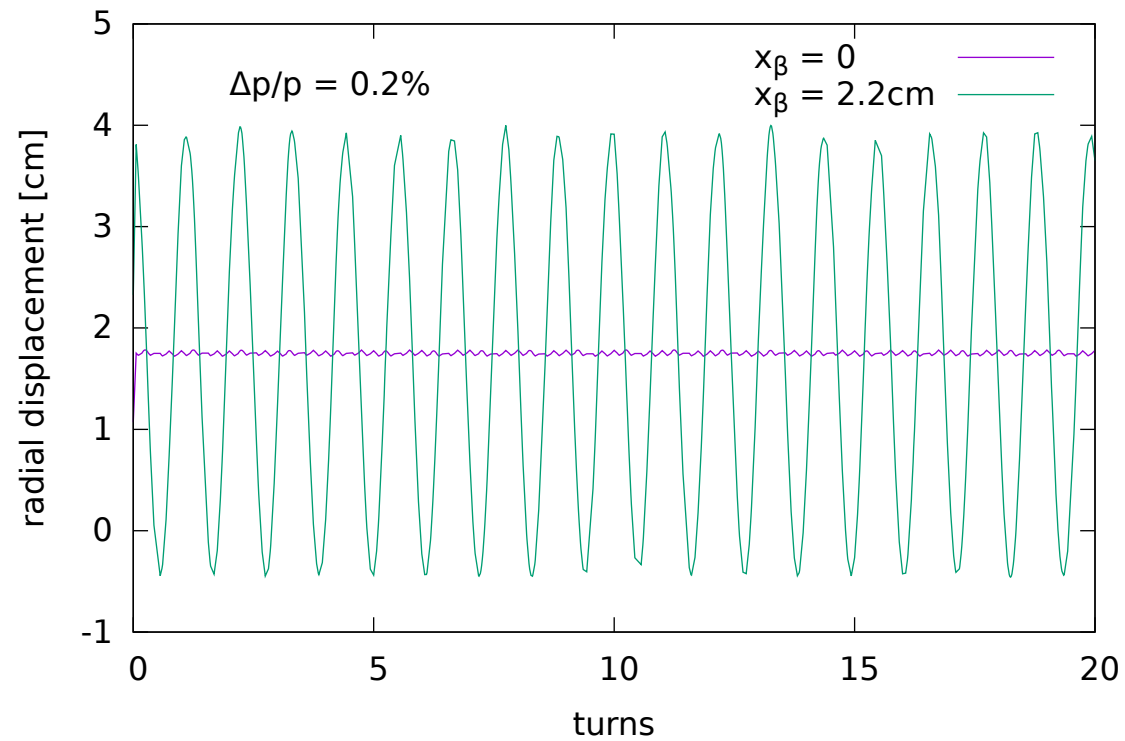
- Multiple points at each momentum correspond to different betatron amplitudes
- The spread at momentum zero is a measure of the accuracy of the simulation (since the E-field correction is nominally zero at the magic momentum)

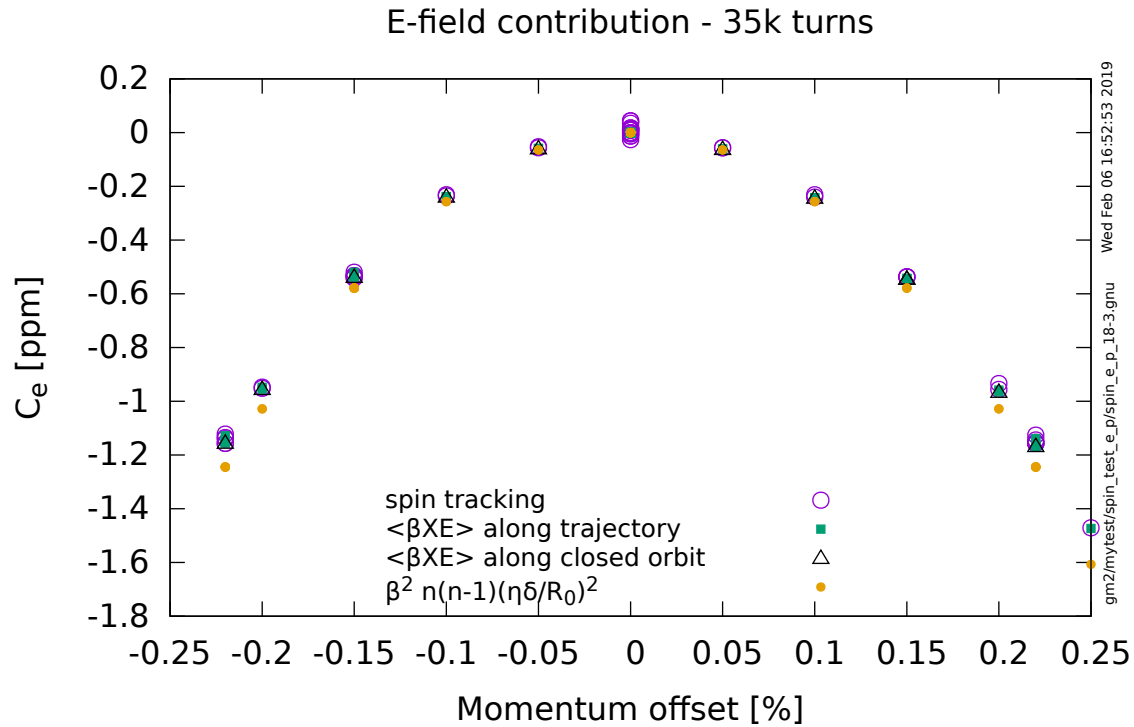


Efield correction computed with spin tracking in good agreement with correction based on $\langle \beta \times \mathbf{E} \rangle$

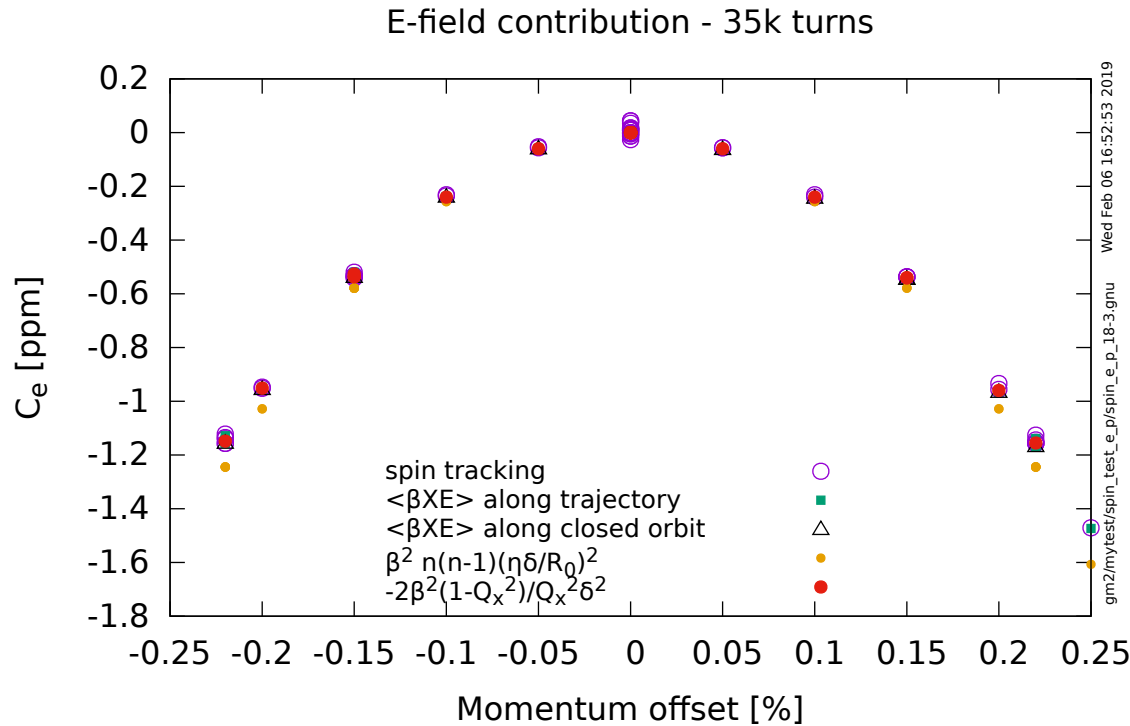


$\langle \beta \times \mathbf{E} \rangle$ along the trajectory is very nearly the same as $\langle \beta \times \mathbf{E} \rangle$ along the closed orbit. (There is little dependence on betatron amplitude)



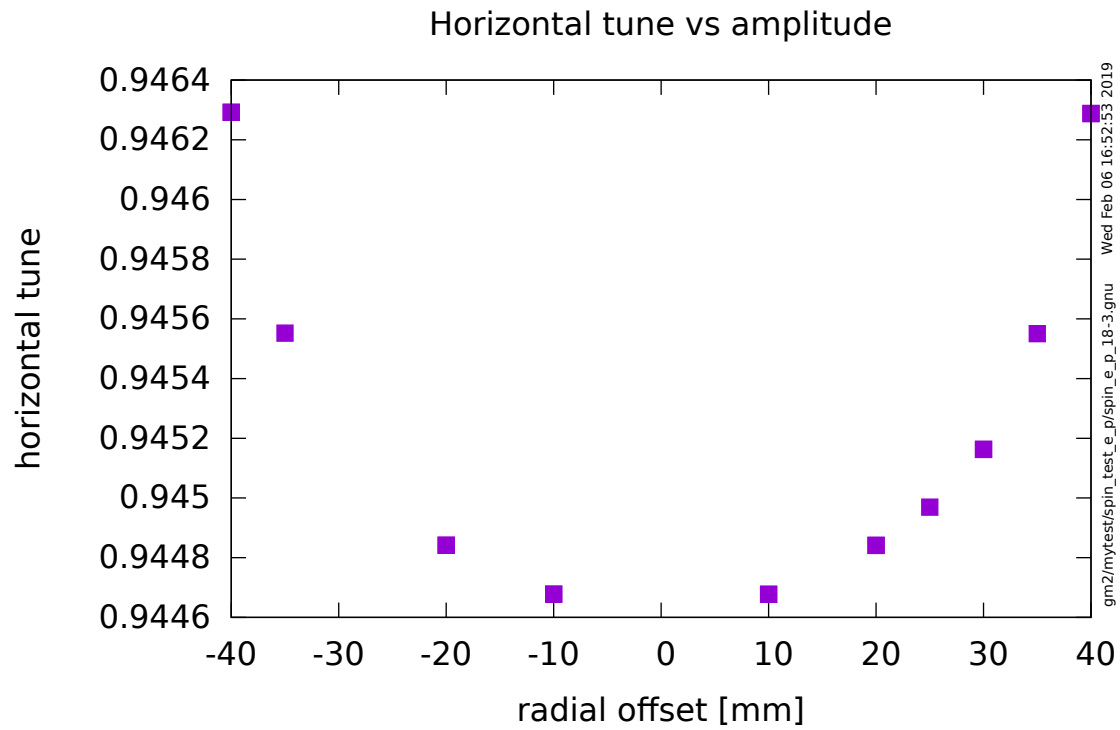


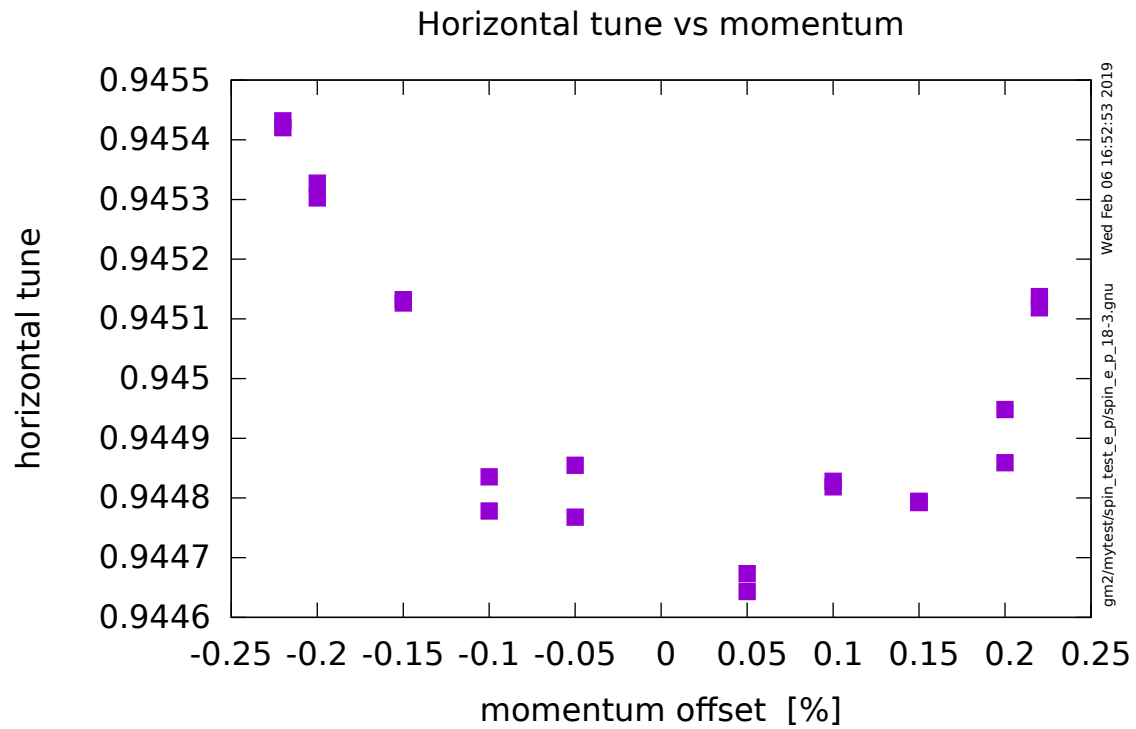
The calculation of the E-field correction that assumes quad linearity, *overestimates* the effect at large momentum offset (where E-field does not increase linearly with displacement)



Replace n , R_0 and η with Q_x and measure Q_x for each momentum $\Rightarrow C_e = -2\beta^2 \frac{(1 - Q_x)^2}{Q_x^2} \delta^2$

The effect of quad nonlinearity can be corrected by measuring momentum dependence of horizontal tune





Comments

- Quad fields are based on an azimuthal slice of 3-D field map with no end effect details
- And perfect relative alignment of plates and absolute alignment about magic radius
- Perfect B-field

Measurement of amplitude/momentum dependence of tunes would be very useful to diagnose quad fields and compensate nonlinearities.